

# Magnetic structural analysis of Pt<sub>3</sub>Fe antiferromagnet by single crystal neutron diffraction

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Pt<sub>3</sub>Fe exhibits antiferromagnetic ordering below  $T_{N1} \sim 170$ K with a propagation wave vector  $\mathbf{q}_1 = (1/2 \ 1/2 \ 0)$ . However, when the chemical composition slightly deviates from the stoichiometric Pt<sub>3</sub>Fe towards Fe rich side, additional magnetic phase transition occurs at  $T_{N2} \sim 110$ K and two types of magnetic ordering with  $\mathbf{q}_1 = (1/2 \ 1/2 \ 0)$  and  $\mathbf{q}_2 = (1/2 \ 0 \ 0)$  coexist [1]. So far, both magnetic domain model and noncollinear magnetic structural model were proposed [2]; however, the magnetic structure is still controversial. In this study, we performed a detailed magnetic structural analysis by a single crystal neutron diffraction to clarify the magnetic structure of non-stoichiometric Pt<sub>3</sub>Fe.

We used the Pt<sub>3</sub>Fe (Fe:25.6 at.%) single crystal with dimensions of  $2.2 \times 2.4 \times 2.6$ mm<sup>3</sup>. Neutron diffraction measurements were performed using four-circle off centered diffractometer (FONDER). We measured approximately 40 and 45 magnetic Bragg reflections for (1/2 1/2 0)-type and (1/2 0 0)-type magnetic structures, respectively, at  $T=10$  K and 120 K.

Figure 1 shows the temperature dependence of the neutron scattering intensity. As the temperature decreases, the (1/2 1/2 0)-type magnetic reflections appear at  $T_{N1}$  and their intensities develop. At  $T=120$  K just above  $T_{N2}$ , the neutron intensity at (0 1/2 1/2), (1/2 0 1/2) and (1/2 1/2 0) is almost the same, indicating that magnetic domains with  $\mathbf{q}_1 = (1/2 \ 1/2 \ 0)$ , (1/2 0 1/2) or (0 1/2 1/2) coexist in a crystal with almost the same volume fraction. As the temperature is further reduced, (1/2 0 0)-type magnetic reflections start to develop below  $T_{N2}$ , associated with the sharp drop of the (1/2 1/2 0)-type ones. However, neutron intensity at (1/2 0 1/2) was not reduced at  $T_{N2}$  and very weakly temperature dependent below  $T_{N2}$ .

To determine the magnetic structure, observed neutron intensities ( $I_{\text{obs}}$ ) were compared with calculated ones ( $I_{\text{cal}}$ ) at  $T=10$  K assuming

magnetic domain model. The  $I_{\text{obs}} - I_{\text{cal}}$  curve shows a good linearity for both types of magnetic structures. At  $T=120$  K above  $T_{N2}$ , the domain model also well explains the observed results and the volume fraction of (0 1/2 1/2), (1/2 0 1/2), (1/2 1/2 0) magnetic structures was 30.4%, 36.8%, 32.8%, respectively. On the other hand, at  $T=10$ K, the volume fraction of (0 1/2 1/2), (1/2 0 1/2), (1/2 1/2 0), (1/2 0 0), (0 1/2 0), (1/2 0 0) magnetic structure was 11.8%, 21.2%, 13.9%, 18.2%, 19.0%, and 16.0%. Such difference indicates that large spin reorientation can occur at  $T_{N2}$  when the (1/2 0 0)-type domains are formed. Noncollinear magnetic structural model might be possible to explain the observed intensities, however, each domain should have different spin tilt angle.

[1] G.E. Bacon et al., Proc. R. Soc. A 272, 387, 1963.

[2] R. Matsui, Y. Tsunoda, J. Phys.: Condens. Matter vol. 21, 124209, 2009.

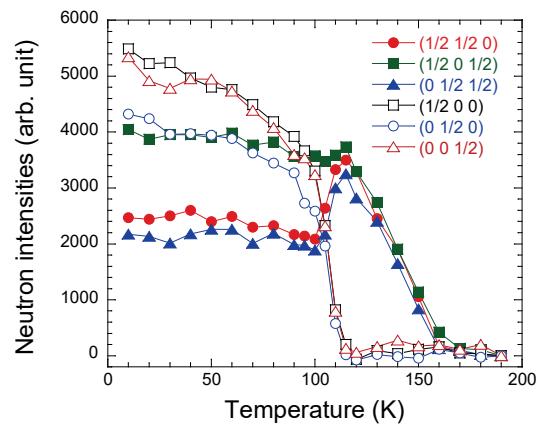


Fig. 1. Temperature dependence of magnetic neutron intensities at (1/2 1/2 0), (1/2 0 1/2), (0 1/2 1/2), (1/2 0 0), (0 1/2 0), (0 0 1/2) magnetic Bragg points.