

Magnetic excitations in a hexagonal $\text{La}_3\text{FeGaS}_7$ with quasi-one-dimensional and chiral crystal structure

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Magnetic systems in which small spin angular momenta $S = 1/2$ or 1 form a one-dimensional chain show quantum nonmagnetic ground states due to large quantum fluctuations. Particularly, in the case of $S = 1$, valence-bond-solid accompanied by Haldane gap realizes as a ground state. For larger spin systems with $S = 3/2, 2 \dots$, on the other hand, classical Neel order manifests itself as the ground state since the quantum fluctuations become small.

Recently, we proposed the one-dimensional magnetic states in hexagonal compounds $\text{La}_3\text{TrGaS}_7$ ($\text{Tr} = \text{Fe}, \text{Co}, \text{Ni}$). The temperature dependences of magnetic susceptibility χ for the $\text{Tr} = \text{Fe}, \text{Co}, \text{Ni}$ systems exhibit a maximum at $T_{\text{max}} = 210, 190, 260$ K, respectively. The maximum suggests developed short-range correlations in the one-dimensional chain. Moreover, the $\chi(T)$ for the $\text{Tr} = \text{Fe}$ system shows an additional maximum and sudden decrease at $T = 20$ and 13 K, respectively, implying long-range magnetic order arising from the finite inter-chain magnetic interaction. On the other hand, the $\chi(T)$ for $\text{Tr} = \text{Ni}$ in 1.8 – 300 K can be reproduced by the calculation of finite Haldane chain system. In this study, we performed powder neutron diffraction and inelastic neutron scattering experiments for $\text{La}_3\text{TrGaS}_7$ ($\text{Tr} = \text{Fe}, \text{Ni}$) in 8 – 200 K at zero magnetic field to extract the magnetic order suggested from the $\chi(T)$ for $\text{Tr} = \text{Fe}$.

Figure 1 shows the powder neutron diffraction patterns for $\text{Tr} = \text{Fe}$ in 8 – 200 K. The pattern at $T = 200$ K can be reproduced well by the simulated pattern with the $\text{La}_3\text{TrGaS}_7$ crystal structure as shown with grey solid line. This pattern is retained from 200 K to 8 K, thus, no additional peak was observed within the measurement accuracy. This result suggests that the $\text{Tr} = \text{Fe}$ system doesn't show neither structural phase transition nor long-range magnetic order. Therefore, the maximum and sudden decrease observed in $\chi(T)$ at $B = 1$ T is

possibly attributed to magnetic field induced phase transition.

Figure 2 shows the inelastic neutron scattering spectra for the $\text{Tr} = \text{Fe}, \text{Ni}$ systems at $q = 2.1$. In the spectra for $\text{Tr} = \text{Ni}$ system, a broad peak was observed around 10 meV, suggesting novel excitation. This excitation energy is comparable to the Haldane gap of 9.2 meV which is estimated from the $\chi(T)$ data. However, similar peak was observed in the $\text{Tr} = \text{Fe}$ system as well. Thus, to confirm whether the excitation in the $\text{Tr} = \text{Ni}$ system is attributed to the Haldane gap or not, it is highly desirable to measure the inelastic neutron scattering spectra for the single-crystalline samples of the $\text{Tr} = \text{Ni}$ system.

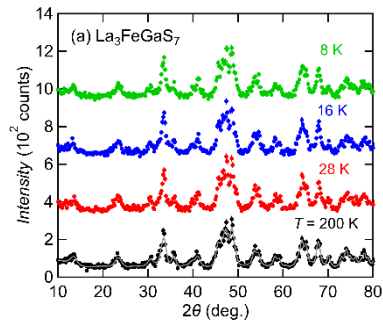


Figure 1. Powder neutron diffraction patterns of $\text{La}_3\text{TrGaS}_7$ $\text{Tr} = \text{Fe}$.

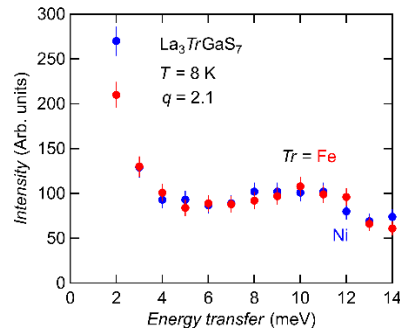


Figure 2. Inelastic neutron scattering spectra for $\text{Tr} = \text{Fe}, \text{Ni}$ at $q = 2.1$ in $T = 8$ K.